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## Effect of Bismuth iodide ( $BiI_3$ ) interfacial layer with different HTL's in FAPI based perovskite solar cell – SCAPS - 1D study

S. Karthick<sup>a, b</sup>, J. Bouclé<sup>b, \*</sup>, S. Velumani<sup>a, c, \*</sup>

(a) Programa de Nanociencias y Nanotecnología, Centro de Investigación y de Estudios Avanzados del Instituto Politécnico Nacional (CINVESTAV-IPN), Av.

Instituto Politécnico Nacional 2508, Col. SanPedro Zacatenco, Ciudad de México, México, Código Postal 07360.

(b) Univ. Limoges, CNRS, XLIM, UMR 7252, F-87000 Limoges, France.

(c) Department of Electrical Engineering (SEES) Centro de Investigación y de Estudios Avanzados del Instituto Politécnico Nacional (CINVESTAV-IPN), Av.

Instituto Politécnico Nacional 2508, Col. SanPedro Zacatenco, Ciudad de México, México, Código Postal 07360.

### Abstract

In this paper, we numerically investigated the effect of Bismuth iodide ( $BiI_3$ ) interfacial layer with different hole transport layer (HTL) candidates (including *Spiro – OMeTAD*,  $Cu_2O$ ,  $CuI$ ,  $CuAlO_2$ ,  $CuSbS_2$ ,  $SrCu_2O_2$ ,  $CuSCN$ , *PTAA*, *P3HT*) in FAPI based perovskite (i.e.,  $FA_{0.85}Cs_{0.15}Pb(I_{0.85}Br_{0.15})_3$ ) solar cells using Solar cell simulator capacitance software (SCAPS-1D). Our results reveal that the addition of a thin  $BiI_3$  layer at the interface between the Perovskite active layer and the HTL efficiently improves hole extraction by defect passivation (i. e., reducing charge recombination and ion migration), which in turn enhances device performance compared to a typical reference architecture. The final optimized device photovoltaic parameters with interfacial layer confirm that the *Cu*-based HTL's, especially  $Cu_2O$  ( $PCE = 24.07\%$ ), and  $SrCu_2O_2$  ( $PCE = 23.91\%$ ) HTL's are more suitable for the  $FA_{0.85}Cs_{0.15}Pb(I_{0.85}Br_{0.15})_3$  solar cells than other HTL alternatives, including *Spiro – OMeTAD* due to higher hole mobility and the valence-band offset alignment between Perovskite/HTL interface. Also, the influence of several metal electrodes (*Ag*, *Cr*, *Cu*, *Au*, *Ni*, *Pt*) is carefully studied with and without the  $BiI_3$  interlayer. It is demonstrated that the energy band misalignment between the HTL and the metallic top electrode restricts charge collection, which is directly associated with low work functions. As a consequence, high work function electrodes such as *Au*, *Ni*, and *Pt* have to be preferred in the presence or absence of the interlayer. The last section addresses the influence of the active layer,  $BiI_3$  interlayer thicknesses on device performance and also the effect of parasitic resistances ( $R_{series}$  and  $R_{shunt}$ ) were studied. From this analysis,  $BiI_3$  interfacial layer seems highly beneficial for improving the performance of experimental perovskite solar cells.

**Keywords:** Bismuth iodide, interfacial layer, perovskite solar cells, SCAPS-1D, HTL.

## 1. Introduction

Due to the rapid improvement in the power conversion efficiency (PCE - up to 25.2 %) of organic-inorganic perovskite solar cells, it gains more attention from the photovoltaic community (NREL, 2020). The intrinsic instability and the presence of lead content in the formamidinium lead iodide (FAPL) perovskites, it forced researchers to develop stabilization (such as cation/halide engineering) and substitution strategies using various elements such as *Sn*, *Bi*, and *Cu*, etc. (Lu et al., 2020). Notably, organic cesium ( $Cs^+$ ) cation and bromide ( $Br^-$ ) halide mixed FAPL perovskite absorber layer (i.e.,  $FA_{1-x}Cs_xPb(I_{1-x}Br_x)_3$ ) shown the remarkable and stable device performance than others because of their slow recombination rate as well as high charge carrier mobility behavior (Groeneveld et al., 2020; McMeekin et al., 2016; Prathapani et al., 2018). Mainly,  $FA_{1-x}Cs_xPb(I_{1-x}Br_x)_3$  combination provides stability enhancement against the moisture and oxygen. Also, it is easier to alter the bandgap value (1.6 – 2 eV) while playing with halide ratio's (iodide and bromide) (Nazarenko et al., 2017; Rehman et al., 2017; Sutter-Fella et al., 2018). These perovskites are more suitable for tandem configurations. Recently, silicon-perovskite based tandem solar cells achieved an excellent efficiency of up to 29 % (Bush et al., 2017; Eperon et al., 2016; Forgács et al., 2017; Köhnen et al., 2019; Lang et al., 2020; NREL, 2020; Zheng et al., 2018).

For efficient photogeneration of electron ( $e^-$ ) and hole ( $h^+$ ) extraction from the perovskite layer and choosing the appropriate  $e^-$  and  $h^+$  transport layer (ETL & HTL) is vital to enhance the PCE. Besides, the interfacial layer (IL) in between the perovskite absorber layer and the HTL interface is considered as a productive layer to avoid the immediate degradation of the absorber layer from external  $H_2O$  &  $O_2$ . Additionally, it helps to reduce the recombination at interfaces to improve the whole device's performance (S. Wu et al., 2019). In general, 64 % of the total solar model cost is related to the HTL (49 %) and the electrode (15 %), the blocking/interlayer cost approximately 0.6% only (Qiu et al., 2019). Besides the price, the organic HTL/electrode (Au or Ag) models suffers from the migration (i.e., metal ions from the metal electrode and halide ions from the perovskite absorber). Usually, there are four possible options to place the interfacial layers based on the planar conventional (n-i-p) and inverted (p-i-n) perovskite structures. Such as in the FTO/ETL interface, ETL/perovskite interface, perovskite/HTL interface, and HTL/electrode interface (Cho and Park, 2017). According to the previous reports, using additional IL in between the perovskite absorber layer and the HTL interface significantly enhances solar cell performance (Cho and Park, 2017; Gharibzadeh et al., 2019; Gil-Escrig et al., 2015; Lee et al., 2014; Y. Liu et al., 2019; Noel et al., 2014; Song et al., 2016). For example, N. K. Noel *et al.* introduced the thiophene and pyridine-based IL between perovskite/HTL interface and demonstrated that the IL improved the device efficiency from 13 % (without IL) to 15.3 % (with thiophene IL) and 16.5 % (with pyridine IL). Because of the passivation strategy, which significantly diminishes the nonradiative recombination pathways (Noel et al., 2014). D. Song *et al.* used the F4TCNQ IL and also achieved an enhanced efficiency from 15 % (without IL) to 18 % (with IL) because of the improved electric field, which mainly decreased the carrier loss at the surface and subsurface of the absorber layer. Also, IL creates an

energy barrier for the photogenerated  $e^-$  which reduces the recombination adjacent to Perovskite/HTL interface (Song et al., 2016). Some research groups used the halide perovskite layer as an IL; for example, J. W. Lee *et al.* deposited the thin  $MAPbI_3$  layer over the  $FAPbI_3$  absorber layer, they got the efficiency enhancement from 15.5% (without IL) to 16 % (with IL) due to an considerable increment in the conduction band minimum at the interface, leading to higher open-circuit voltage (Lee et al., 2014). Similarly, K. T. Cho *et al.* achieved an increased PCE from 18.9 % (without IL) to 21.3 % (with IL) using additional  $FAPbBr_{3-x}I_x$  layer over the  $(FAPbI_3)_{0.85}(MAPbBr_3)_{0.15}$  perovskite absorber, which efficiently suppresses the interfacial charge recombinations (Cho et al., 2017). In this perspective, using bismuth (*Bi*) based IL in perovskite device will be beneficial for the device performance. For example, L. Fu *et al.* employed the bismuth telluride ( $Bi_2Te_3$ ) IL in HTL/Inorganic-Perovskite interface, and achieved an improved device performance from 7.4 % (without) to 11.9 % (with IL). These results clearly explain that, IL efficiently blocks the iodide ions migrations from the absorber layer (Fu et al., 2019). Also, bismuth iodide ( $BiI_3$ ) material gains more attention due to its non-toxic behavior, a narrow bandgap (1.72 eV), a high absorption coefficient, and short carrier lifetimes, considering these properties of  $BiI_3$ , it is more suitable for photovoltaic applications (Brandt et al., 2015a; Coutinho et al., 2019; Hamdeh et al., 2016; Hu et al., 2020; Lee et al., 2018; Yoo et al., 2019). Moreover, according to the U. S. national mineral information center report, the Bismuth price is relatively too low compared to the lead (National Minerals Information Center, n.d.). Recently, Y. Hu *et al.* deposited the pristine  $BiI_3$  IL over the titanium oxide ETL (i.e., ETL/Perovskite interface), and found that it improves the electron transport behavior by passivating the interfacial trap states. As a result, they got an increased PCE from 13.8 % (without) to 17.8 % (with IL) (Hu et al., 2020). But still, there is no  $BiI_3$  IL-based simulation/theoretical analysis reported so far in the Perovskite/HTL interface.

Therefore, in this context, we theoretically introduce a thin  $BiI_3$  IL between  $FA_{0.85}Cs_{0.15}Pb(I_{0.85}Br_{0.15})_3$  absorber/Spiro-OMeTAD (HTL) layer using solar cell simulator capacitance software (SCAPS-1D) to understand the beneficial effect of  $BiI_3$  on the device performance. Firstly, we construct an ideal n-i-p planar device without using the IL layer as a reference, and after placing a thin  $BiI_3$  IL, our simulated device (FTO/ETL/Perovskite/ $BiI_3$ /HTL/Electrode) showed enhanced performance compared to the referenced one. Secondly, we carefully investigated the device performance using different HTL candidates, such as 2,2',7,7'-Tetrakis[N, N-di(4-methoxyphenyl)amino]-9,9'-spirobifluorene (*Spiro* – *OMeTAD*), cuprous oxide ( $Cu_2O$ ), Strontium Cuprate ( $SrCu_2O_2$ ), Cuprous Aluminate ( $CuAlO_2$ ), copper antimony sulfide ( $CuSbS_2$ ), cupric thiocyanate ( $CuSCN$ ), Cuprous iodide ( $CuI$ ), Poly(triarylamine) (*PTAA*), poly(3-hexylthiophene-2,5-diyl) (*P3HT*) with and without  $BiI_3$  IL. The final optimized device photovoltaic parameters with IL confirms that the *Cu*-based  $Cu_2O$ , and  $SrCu_2O_2$  HTL's are more suitable than other HTL nominees due to their band alignment with  $BiI_3$  IL. Furthermore, we evaluated the effect of several electrodes (*Ag, Cr, Cu, Au, Ni, Pt*) based on their work function (WF) value with the above mentioned HTL's. At last, we studied the influence of IL thickness over the perovskite layer thickness and the effect of parasitic resistances ( $R_{series}$  and  $R_{shunt}$ ) with suitable HTL. Our results deliver insightful facts to understand the effect of IL over the planar n-i-p *FAPbI* based perovskite device

structure, which is useful to enhance the PCE. As a result, using  $BiI_3$  IL in between the perovskite/HTL interface could help to improve the experimental device performance.

## 2. Ideal device simulation methodology

Our simulated n-i-p device composed of  $FTO$  (transparent conductive oxide)/ $SnO_2$  (ETL)/ $FA_{0.85}Cs_{0.15}Pb(I_{0.85}Br_{0.15})_3$  (perovskite absorber layer)/ $BiI_3$  (with and without - IL)/ $Spiro - OMeTAD$  (HTL) / $Au$  (electrode), shown in **Fig. 1**. In this study, all the simulations were carried out using solar cell simulator capacitance software (SCAPS-1D, version 3.3.07) to solve the Poisson and continuity equations (**Burgelman et al., 2000**). The simulated device was illuminated under one sun AM 1.5 G ( $100 \text{ mW cm}^{-2}$ ) solar spectrum with 300 K temperature, and the parasitic resistances ( $R_{Series}$  and  $R_{Shunt}$ ) are not taken into these initial simulations. After illumination (photons strike) from ETL to HTL side (see **Fig.1**), the perovskite layer creates charge carriers, and diffuse into the electron ( $e^-$ ) and hole ( $h^+$ ) transport layers. Finally,  $e^-$ 's and  $h^+$ 's were collected by desired electrodes. The optimized absorber layer (500 nm), ETL (70 nm), all the HTL (150 nm), and IL (30 nm) thicknesses, as well as all other parameters such as  $e^-$  and  $h^+$  mobility,  $e^-$  affinity, permittivity, the effective density of states, doping densities, and defect densities, are taken from previously published experimental and theoretical results, listed in Table I. (**Abdelaziz et al., 2020; Azri et al., 2019; Brandt et al., 2015b; Coulibaly et al., 2019; Coutinho et al., 2019; Han et al., 2014; Lee et al., 2018; Lin et al., 2020; Prathapani et al., 2018; Shasti and Mortezaali, 2019**) The employed  $e^-$  and  $h^+$  velocity values are considered as  $1 \times 10^7 \text{ cm/s}$ . The work function of TCO (front contact - FTO) and electrode (back contact - Au) are set to 4.4 eV and 5.1 eV, respectively. All the simulation was done, without considering the optical reflectance of each layers at the interface nor surface. Simulating ideal devices is a first step to emphasize the influence of the IL on device operation, independently of the main electrical losses that are found in realistic devices. It is designed to help the experimentalist to tailor their experiments (**Karthick et al., 2020**). The effect of parasitic resistances ( $R_{Series}$  &  $R_{Shunt}$ ) were demonstrated at the end of this paper. The energy band alignment between the absorber layer, ETL, IL, and all HTL illustrated in **Fig. 2**.

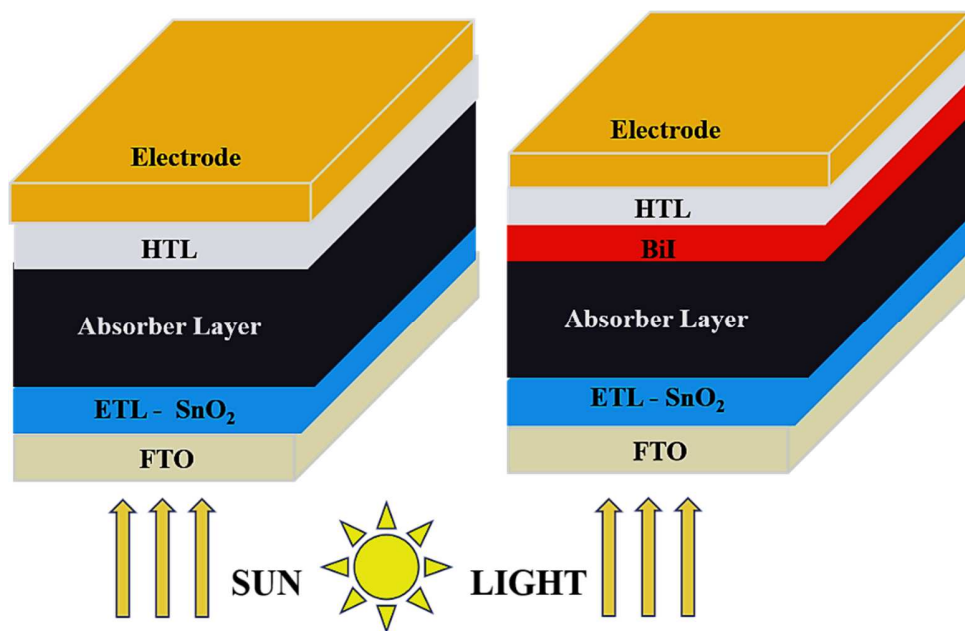


Fig. 1 The n-i-p typical perovskite solar cell structure (left) and with  $BiI_3$  IL (right).

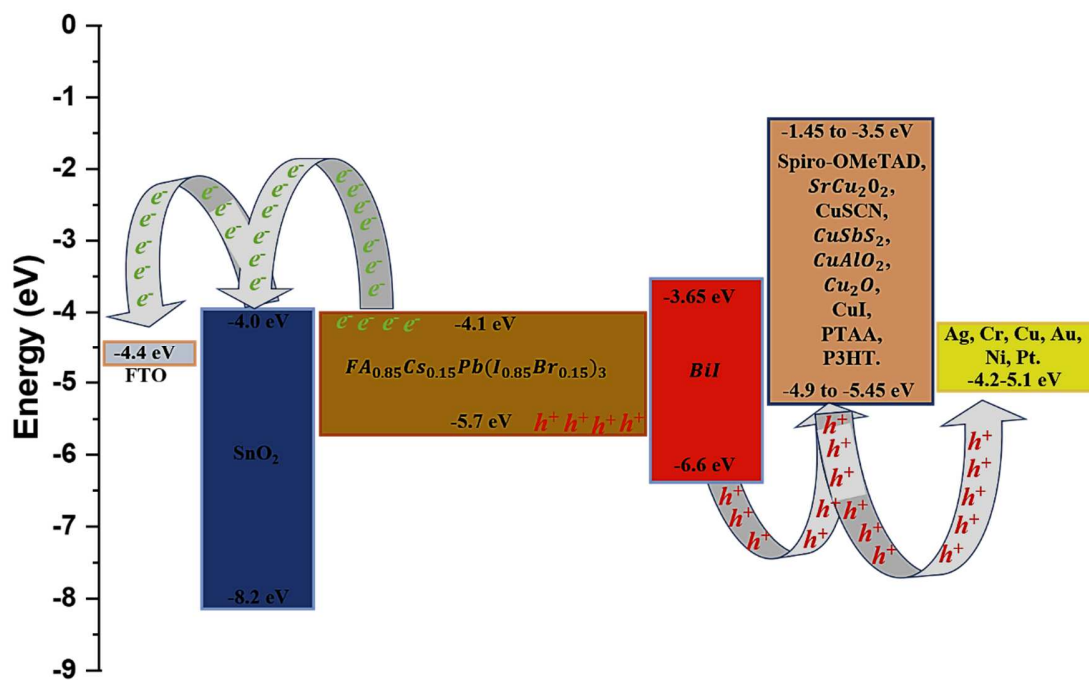


Fig. 2 Band alignments between ETL, different HTL's with perovskite absorber including  $BiI_3$  IL.

Table I: Details of primary input parameters used for the simulation of *FAP*I based perovskite solar cells.

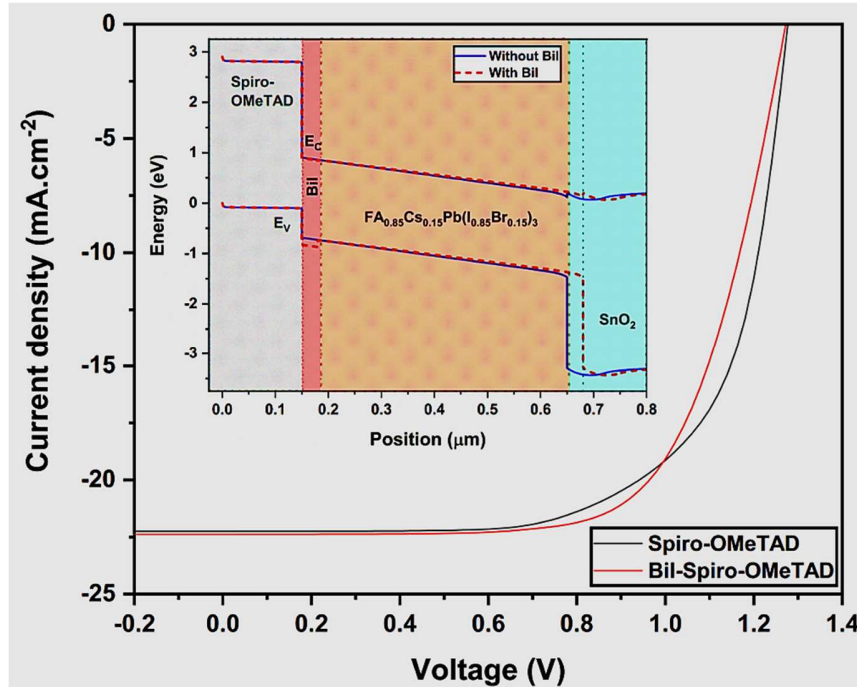
Parameters	FTO	ETL (SnO <sub>2</sub> )	FA <sub>0.85</sub> CS <sub>0.15</sub> Pb (I <sub>0.85</sub> Br <sub>0.15</sub> ) <sub>3</sub>	BiI <sub>3</sub> (IL)	HTL (Spiro-OMETAD)
<i>Thickness (nm)</i>	500	70 (Karthick et al., 2020)	500	30	150
<i>Band gap (eV)</i>	3.50	3.50	1.59	1.72 (Coutinho et al., 2019)	2.9
<i>E Affinity</i>	4.00	4.00	4.09 (Prathapani et al., 2018)	4.10 (Lee et al., 2018)	2.2
<i>Permittivity</i>	9.00	9.00	6.600	5.78 (Brandt et al., 2015b)	3.00
<i>Effective density of states at CB</i>	2.2×10 <sup>18</sup>	2.2×10 <sup>17</sup>	2.0×10 <sup>19</sup>	2.5×10 <sup>19</sup>	2.2×10 <sup>19</sup>
<i>Effective density of states at VB</i>	2.2×10 <sup>18</sup>	2.2×10 <sup>17</sup>	2.0×10 <sup>18</sup>	2.5×10 <sup>19</sup>	2.2×10 <sup>19</sup>
<i>Mobility of e<sup>-</sup></i>	20	20	8.16 (Prathapani et al., 2018)	600 (Han et al., 2014)	1.0×10 <sup>-4</sup>
<i>Mobility of h<sup>+</sup></i>	10	10	2	200 (Callahan et al., 2018)	1.0×10 <sup>-4</sup>
<i>Density of n-type doping</i>	1.0×10 <sup>15</sup>	1.0×10 <sup>17</sup>	1.3×10 <sup>16</sup>	1.0×10 <sup>16</sup>	0
<i>Density of p-type doping</i>	0	0	1.3×10 <sup>16</sup>	1.0×10 <sup>16</sup>	1.0×10 <sup>18</sup>
<i>Density of defects</i>	Donor- 1.0×10 <sup>18</sup>	Donor – 1.0×10 <sup>15</sup>	Neutral- 4×10 <sup>13</sup>	Neutral- 1×10 <sup>15</sup>	Acceptor- 1.0×10 <sup>15</sup>

### 3. Results and discussion

#### 3.1 Ideal device performance

In this simulation, Firstly SnO<sub>2</sub> (ETL) and Spiro – OMeTAD (HTL) based *FAP*I perovskite ideal device performances are carefully investigated as a reference, which shows a PCE of 19.17 % with a short current density ( $J_{sc}$ ) value of 22.24 mA/cm<sup>-2</sup>, a fill factor (FF) of 67.46 %, and an open-circuit voltage ( $V_{oc}$ ) of 1.27 V that is associated with previously reported results (McMeekin et al., 2017; Yu et al., 2016). Here, we used the SnO<sub>2</sub> as an ETL instead of TiO<sub>2</sub>

and  $ZnO$  due to the excellent properties of  $SnO_2$  (Jiang et al., 2018; D. Liu et al., 2019) and also,  $SnO_2$  provides perfect band alignment with mixed cation/halide perovskite than  $TiO_2$  (Correa Baena et al., 2015).



**Fig. 3** Current density- voltage (J-V) characteristic curves of both ideal device structures (with and without  $Bil_3$  IL). The inset represents the energy band diagram.

The primary and crucial job of the ETL layer is to extract the electron from the absorber layer and simultaneously blocking the holes to avoid recombinations. Experimentally, due to the different ETL layer thicknesses, the light transmittance of the film will vary. Therefore, the ETL layer thickness plays a vital role (Xiong et al., 2018). In general,  $SnO_2$  – ETL layer needs a higher thickness than  $TiO_2$  – ETL to completely cover the TCO substrate due to the larger size of  $SnO_2$  nanoparticles (D. Liu et al., 2019). Fumin Li et al. studies demonstrate that the 66 nm  $SnO_2$  – ETL device showing an improved PV performance than 34 nm and 48 nm ETL devices. The thicker ETL effectively blocks the hole transport, and it reduces the charge carrier recombination at the interfaces, which is beneficial for the device's performance (Li et al., 2018). Also, the 70 nm  $Cl - SnO_2$  – ETL device provides an enhanced PCE than 40 nm and 90 nm layer contained devices due to tight interface contact with the TCO substrate and absorber (Duan et al., 2017).

Furthermore, it is well-known that several factors affect perovskite device performance, including HTL layer thickness. In this simulation, we used an optimized thickness of 150 nm for all different HTL layers. Previous experimental reports justify that the 180 nm HTL layer device offers a higher PV performance than other HTL thicknesses, such as 100, 250, 400, 450, 500, and 700 nm (Kim et al., 2015). Apart from the hole extraction and transportation, the absorber layer's safety and suppression of the interface defects are crucial to achieving good device performance. Therefore, the HTL thickness should be thicker than 150 nm to avoid direct contact between the



absorber and the electrode. Also, a thick HTL efficiently enhances the charge carrier collection due to an improved reflection from the smooth HTL/electrode interface; thus, the device efficiency is highly improved (Lei et al., 2018; Marinova et al., 2015).

Later, we introduced an additional thin  $BiI_3$  IL in between the perovskite/HTL interface (see **Fig.1**). Perovskite/HTL interface engineering is a well-known strategy to efficiently passivate the defects, which generally promotes hole extraction by blocking the secondary electrons (**Gil-Escrig et al., 2015**). Surprisingly, here all the photovoltaic parameters are drastically enhanced, the obtained PCE improved from 19.17 % to 19.28 % ( $J_{sc} = 22.38 \text{ mA/cm}^{-2}$ ,  $FF = 67.73 \%$ , and  $V_{oc} = 1.27$ ). **Fig. 3** illustrates the simulated current density – voltage curves (J-V) of both ideal device structures, and the inset's shows the energy band diagram. The obtained quantum efficiency (QE) graphs were displayed in the supplementary information (see **Fig. S1**). Experimentally, introducing IL is not only beneficial for stopping ion migration, and also it modifies the energy level alignment between the layers. In general, the interfacial energy level alignment influence the hole injection rate (i.e., perovskite/HTL interface). The energy barrier (EB) at the interfaces leads to charge carrier recombination losses restricting the charge transfer. The non-appearance of EB at the interfaces helps the smooth and fast charge transfer/injection and efficiently reduces the recombinations (Shao and Loi, 2020). In this case,  $BiI_3$  has a higher bandgap value and larger conduction band than the absorber (see **Table I and Fig. 2**), leading to the favored energy level alignment that efficiently minimizes/reduces the EB at the HTL/perovskite interface.

Henceforth, introducing  $BiI_3$  IL in the Perovskite/HTL interface facilitates the interfacial hole transport and possibly suppresses the interfacial charge recombination. The valence band offset (VBO) is generally defined as the difference between the  $e^-$  affinity and bandgap of the perovskite absorber; and  $e^-$  affinity and bandgap of the HTL (i.e.,  $VBO = e^- \text{ affinity}_{HTL} - e^- \text{ affinity}_{perovskite} + bandgap_{HTL} - bandgap_{Perovskite}$ ) (**Abdelaziz et al., 2020; Sahu and Dixit, 2018; Shasti and Mortezaali, 2019**). The negative VBO is obtained in the perovskite/*Spiro-OMeTAD* interface (-0.58 eV), which signifies that the presence of recombination of minority charge carriers occurs in the interfaces (**Lin et al., 2019**). After adding the  $BiI_3$  IL, the VBO (i.e.,  $VBO = e^- \text{ affinity}_{BiI_3} - e^- \text{ affinity}_{HTL} + bandgap_{BiI_3} - bandgap_{HTL}$ ) (**Minemoto and Murata, 2015**) turns positive (0.72 eV), by means that the band alignment is modified between the absorber and the HTL interface, which is beneficial for balanced charge carrier transfer; as a result, it boosts the FF value (see **Table S1**). Additionally, the current value of  $BiI_3$  IL attached device is higher than that of the reference device, which generally leads to higher charge carrier mobility, helping further PCE improvement.

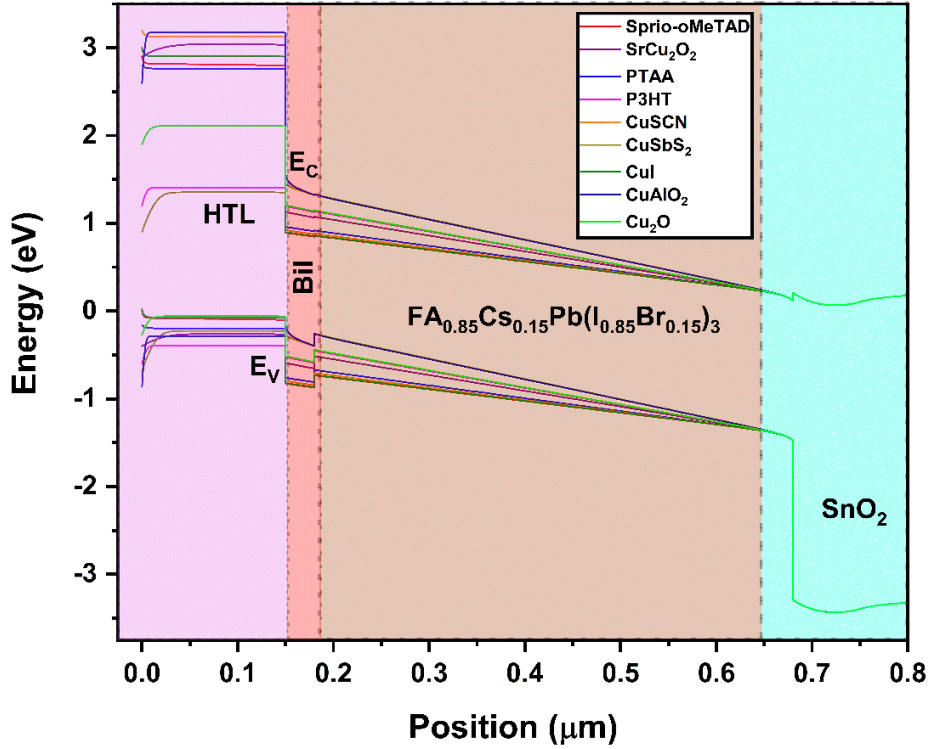
Hence, it is clear that the  $BiI_3$  IL influences the device results, so in the following sections, we are going to check the device performances based on several HTL candidates with and without  $BiI_3$  IL to understand the interfacial layer effect. Finally, we will correlate the influence of the  $BiI_3$  IL thickness with the perovskite absorber layer thickness and also the effect of parasitic resistances, using the most appropriate HTL system.

### 3.2 Effect of different HTL's with and without $BiI_3$ IL

Experimentally, *Spiro – OMeTAD* HTL needs some external doping (such as tri-tert-butylpyridine (tBp), lithium bis-tri-fluoro-methane-sulfonimide (Li-TFSI), and other cobalt salts) to improve the internal charge mobility, but at the same time, it is rising the interface recombination, which led to poor device performance (**Gheno et al., 2016**). Therefore, finding alternative low-cost and stable HTL is essential to improve device efficiency. Hence, in this section, several HTL ( $Cu_2O$ ,  $CuI$ ,  $CuAlO_2$ ,  $CuSbS_2$ ,  $SrCu_2O_2$ ,  $CuSCN$ ,  $PTAA$ ,  $P3HT$ ) candidates are studied and presented. Input parameters of the mentioned HTL materials are listed in Table II. By using 70 nm of  $SnO_2$  as an ETL and switching different HTL materials with and without  $BiI_3$  IL, the comparative energy band diagram is illustrated in **Fig. 4** and **Fig. S2**. Obtained J-V characteristic curves, as well as the appropriate energy band diagrams (inset) for all mentioned HTL devices, are presented in **Fig. 5**, **Fig. 6**, and **Fig. S3 - Fig. S8**, respectively.

Table II: Input parameters for proposed different HTL materials.

Parameters	$SrCu_2O_2$	$Cu_2O$	$CuAlO_2$	$CuSCN$	$CuSbS_2$	$CuI$	$PTAA$	$P3HT$
<i>Thickness (nm)</i>	150	150	150	150	150	150	150	150
<i>Bandgap (eV)</i>	3.3	2.17	3.46	3.2	1.58	2.98	2.96	1.8
<i>E Affinity</i>	2.2	3.2	2.5	1.9	4.2	2.1	2.3	3.9
<i>Permittivity</i>	9.77	7.1	60	10	14.6	6.5	9	3
<i>Effective density of states at CB</i>	$2.0 \times 10^{20}$	$2.1 \times 10^{17}$	$2.0 \times 10^{20}$	$2.2 \times 10^{19}$	$2.2 \times 10^{18}$	$2.8 \times 10^{19}$	$2.0 \times 10^{21}$	$2.0 \times 10^{20}$
<i>Effective density of states at VB</i>	$2.0 \times 10^{21}$	$1.0 \times 10^{19}$	$1.0 \times 10^{22}$	$1.8 \times 10^{19}$	$2.2 \times 10^{19}$	$2.0 \times 10^{19}$	$2.0 \times 10^{21}$	$2.0 \times 10^{20}$
<i>Mobility of <math>e^-</math></i>	0.1	200	2	$1.0 \times 10^{-4}$	49	100	1	$1.0 \times 10^{-4}$
<i>Mobility of <math>h^+</math></i>	0.46	80	8.6	$1.0 \times 10^{-1}$	49	43.9	40	$1.0 \times 10^{-3}$
<i>Density of n-type doping</i>	0	0	0	0	0	0	0	0
<i>Density of p-type doping</i>	$1.0 \times 10^{17}$	$1.0 \times 10^{18}$	$1.0 \times 10^{20}$	$1.0 \times 10^{18}$	$1.0 \times 10^{18}$	$1.0 \times 10^{18}$	$1.0 \times 10^{18}$	$1.0 \times 10^{18}$
<i>Density of defects</i>	Acceptor- $1.0 \times 10^{15}$	Acceptor- $1.0 \times 10^{15}$	Acceptor- $1.0 \times 10^{15}$	Acceptor- $1.0 \times 10^{15}$	Acceptor- $1.0 \times 10^{15}$	Acceptor- $1.0 \times 10^{15}$	Acceptor- $1.0 \times 10^{15}$	Acceptor- $1.0 \times 10^{15}$
<i>References</i>	(Shasti and Mortezaali, 2019)	(Shasti and Mortezaali, 2019) (Lin et al., 2020)	(Shasti and Mortezaali, 2019)	(Azri et al., 2019) (Chakraborty et al., 2019)	(Devi and Mehra, 2019)	(Coulibaly et al., 2019) (Azri et al., 2019)	(Coulibaly et al., 2019)	(Coulibaly et al., 2019) (Azri et al., 2019)

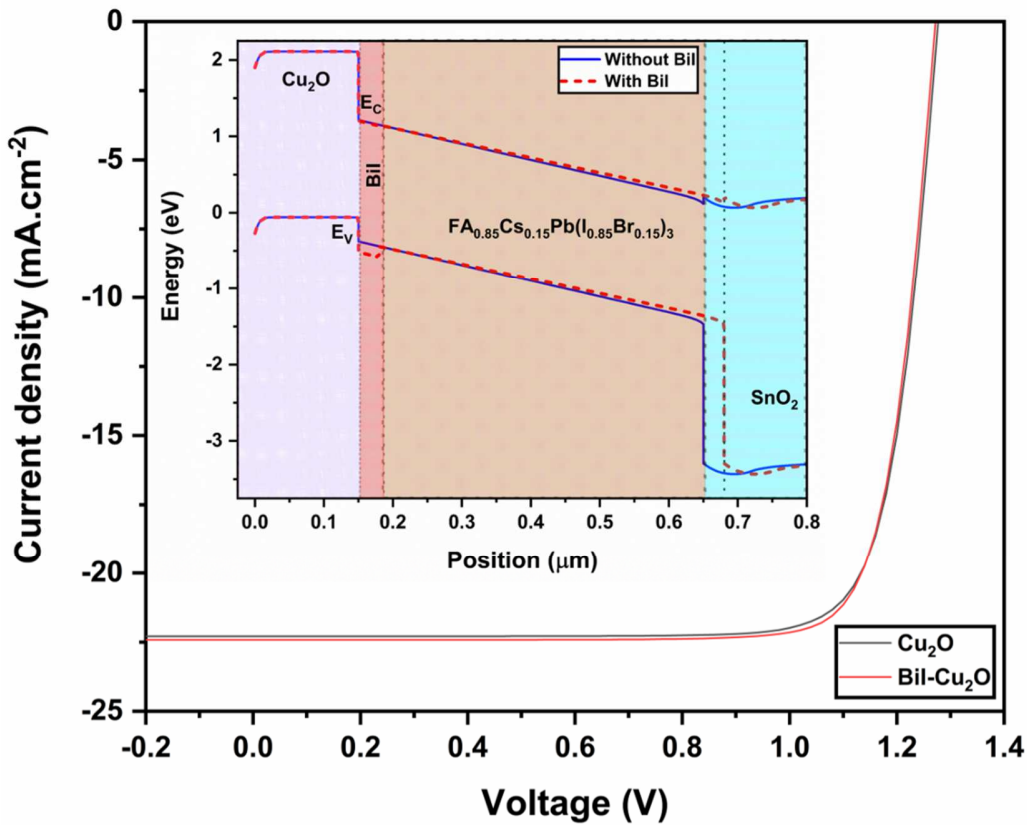


**Fig. 4** Energy band diagram for the reference device including  $BiI_3$  IL with different HTL candidates.

After testing several HTL candidates over the same configured structure, we noticed that the  $Cu_2O$ -HTL based device shows a remarkable PCE of 23.07 % than other competitors due to their excellent hole mobility value (Table II) and the energy band alignment with the absorber layer (see Fig. 5). The obtained FF value (81%) is highly improved compared to the reference *Spiro* – *OMeTAD*-HTL device (67 %), which can be attributed to the promotion of hole conduction through the HTL. The simulated photovoltaic parameters of all different HTL based device results are illustrated in Table III. The VBO is significantly improved (- 0.31) while using  $Cu_2O$ -HTL instead of *Spiro* – *OMeTAD*, that could be a good reason for the enhanced device performance. After adding the  $BiI_3$  IL in between the Perovskite/ $Cu_2O$ -HTL interface, as we expected, the overall device performance increased from 23.07 % to 23.29 % compared to the standard device (see Fig. 5 and Table III) due to the positive VBO (0.45 eV). The VBO value of  $BiI_3/Cu_2O$  interface is lower than the  $BiI_3/Spiro$  – *OMeTAD* interface, yet the device performance is boosted because of the high VBO, which acts as a barrier for the hole diffusion (Minemoto and Murata, 2015) and lead to lower performance (i.e., *Spiro* – *OMeTAD*-HTL device). Also, hole mobility plays a massive role in efficiently extracting and surpassing the holes from the absorber to HTL and the electrode (Table II). For example, in a real device, Elesman *et al.* experimental TRPL results explain the  $Cu_2O$ /perovskite films exhibit excellent lifetime than P3HT/perovskite films, which signifies that the  $Cu_2O$ /perovskite interface gains more hole extraction than P3HT/perovskite interface (Elesman et al., 2019). The  $SrCu_2O_2$  – HTL device having the second-highest efficiency of 22.75 %, here also the VBO is negative (-

0.18 eV) but it is better than the *Spiro – OMeTAD* and *Cu<sub>2</sub>O*-HTL's, at the same time, the hole mobility is lower than the *Cu<sub>2</sub>O* and higher than *Spiro – OMeTAD* (Table II). The addition of *BiI<sub>3</sub>* IL furthermore improves the VBO to 0.32 eV, which enhances the device performance from 22.75 % to 23.20 % (see Fig. 6), it is higher than the previously simulated perovskite device results (Table III). The *PTAA*-HTL based device showed PCE of 19.80 %, which is attributed to the VBO (- 0.42 eV) and hole mobility value, it is finer than *Spiro – OMeTAD*-HTL. While adding the *BiI<sub>3</sub>* IL, the VBO is considerably enhanced (0.56 eV), which directed towards the higher performance from 19.80 % to 20.15 % due to an enhancement in the interface (see Fig. S3). The previously reported simulated *PTAA*-HTL device shows higher efficiency (23.5 %) than the current device results because of the higher HTL thickness (1000 nm) and the  $J_{sc}$  value (41.03 mA/cm<sup>-2</sup>) (Coulibaly et al., 2019). *CuAlO<sub>2</sub>*-HTL based device displays 19.98 % of efficiency due to the reduced charge transfer resistance between HTL and the electrode (Igbari et al., 2015), the obtained performance is higher than the reference *Spiro – OMeTAD*-HTL device (Table III). In this case, the VBO is positive (0.28 eV) compared to formerly discussed HTL contenders; surprisingly, after including *BiI<sub>3</sub>* IL, the VBO turns into negative (- 0.14 eV), as a consequence, the  $V_{oc}$  value is significantly reduced. Even though, the performance is slightly increased from 19.98 % to 20.04 % due to an enrichment in short-circuit current (see Fig. S4, and Table III). *CuSbS<sub>2</sub>*-HTL device demonstrating the PCE of 19.08 %, here also the VBO is positive (0.1 eV), while adding *BiI<sub>3</sub>* IL, the VBO value is reduced a little bit (0.04 eV), which is reflecting in the FF result; however, the PCE is slightly improved (see Fig. S5, and Table III). Previous SCAPS reports shows that a higher efficiency was obtained using *CuSbS<sub>2</sub>*-HTL (24.1 %) because of the larger  $J_{sc}$  value (31.7 mA/cm<sup>-2</sup>) and the HTL thickness (438 nm), (Devi and Mehra, 2019) is much larger than the current result (22.6 mA/cm<sup>-2</sup>), shown in Table III. *CuSCN*-HTL device displaying PCE of 19.05 %, the VBO is negative (-0.58 eV), after adding *BiI<sub>3</sub>* IL, the VBO is significantly improved to 0.72 eV, at the end, the FF is reduced due to the higher VBO, but the device shows the same performance because of the improvement in the  $J_{sc}$  value (see Fig. S6, and Table III). Previous report shows that the device performance is not entirely associated with perovskite/*CuSCN* interface rather than *CuSCN*/electrode (Arora et al., 2017). Therefore, we will demonstrate the effect of HTL with different electrodes in the following sections. But in reality, experimental *CuSCN*-HTL contained device demonstrates excellent thermal degradation stability than well known *Spiro – OMeTAD*-HTL due to the strong resistance against environmental stress (Jung et al., 2016). *CuI*-HTL device exhibiting PCE of 18.81 %, the VBO is negative (-0.6 eV), after inserting *BiI<sub>3</sub>* IL, the VBO is drastically increased (0.74 eV), the performance is also improved minorly but not more (i.e., 18.81 % to 18.86 %) (see Fig. S7 and Table III), the high VBO might restrict the further improvement as we discussed before. Noticeably, *P3HT*-HTL device showing the lowest photovoltaic performance (17.99 %) among all HTL's due to the lower hole mobility value and the lower VBO (0.02 eV). According to the previous report, the significant mismatch between the quasi-Fermi level splitting (QFLS) to  $V_{oc}$  affects the results that occur while using the *P3HT*-HTL compared to *PTAA*-HTL, enhancing the interfacial recombination at the interfaces (Caprioglio et al., 2019) might be another reason for the lower performance. After including the *BiI<sub>3</sub>* IL, the VBO is changed (0.12 eV), but surprisingly, the FF and  $V_{oc}$  values are dropped, it further reduces the performance (see Fig. S8, Table III). Overall, other HTL based device exhibits lower performance than these two

( $Cu_2O$  &  $SrCu_2O_2$ ) with and without  $BiI_3$  IL. According to the previous report, the vacancy formation energies are reduced in the order of  $Cu_2O < SrCu_2O_2 < CuAlO_2$  and it affects the conduction properties; mainly, it modifies the acceptor defects. The vacancy distribution in  $Cu_2O$  impacts the bandgap, and the conduction properties (Nolan, 2008). Our simulated device results are compared with recently published reports, shown in Table III. It emphasizes that the simulated device performance is solely based on the appropriate absorber layer, ETL, HTL arrangement, and energy band alignment (Azri et al., 2019; Coulibaly et al., 2019; Elseman et al., 2019; Lin et al., 2020; Shasti and Mortezaali, 2019; Zheng et al., 2017). Therefore, the simulated device with  $BiI_3$  IL efficiently boosting up most of the above-mentioned HTL's based device performances compared to the referenced device, especially, the  $J_{sc}$  and FF values are improved (Table III). The obtained QE behavior of *Spiro* – *OMeTAD*-HTL referenced device is nearly similar to the  $BiI_3$  IL contained device, shown in Fig. S1. Hence, it is an evident that the  $BiI_3$  IL with suitable HTL ( $Cu_2O$  &  $SrCu_2O_2$ ) directly influences the device performance. In the following sections, we will examine the various electrodes with and without  $BiI_3$  IL (using different HTL's) to understand the effect over the HTL/electrode interface.



**Fig. 5** shows the Current density- voltage (J-V) characteristic curves of both ideal device structures (with and without  $BiI_3$  IL) with  $Cu_2O$ -HTL. The insets represent the energy band diagram.

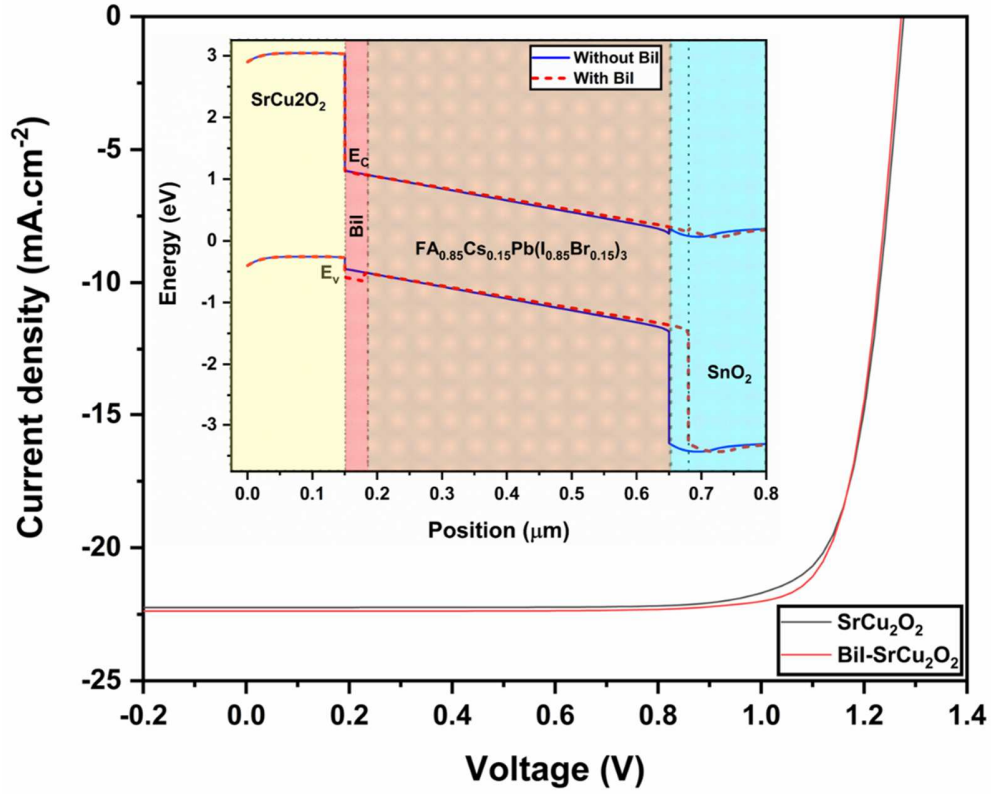


Fig. 6 shows the Current density- voltage (J-V) characteristic curves of both ideal device structures (with and without  $BiI_3$  IL) with  $SrCu_2O_2$  – HTL. The insets represent the energy band diagram.

Table III: Recorded photovoltaic parameters for proposed different HTL materials.

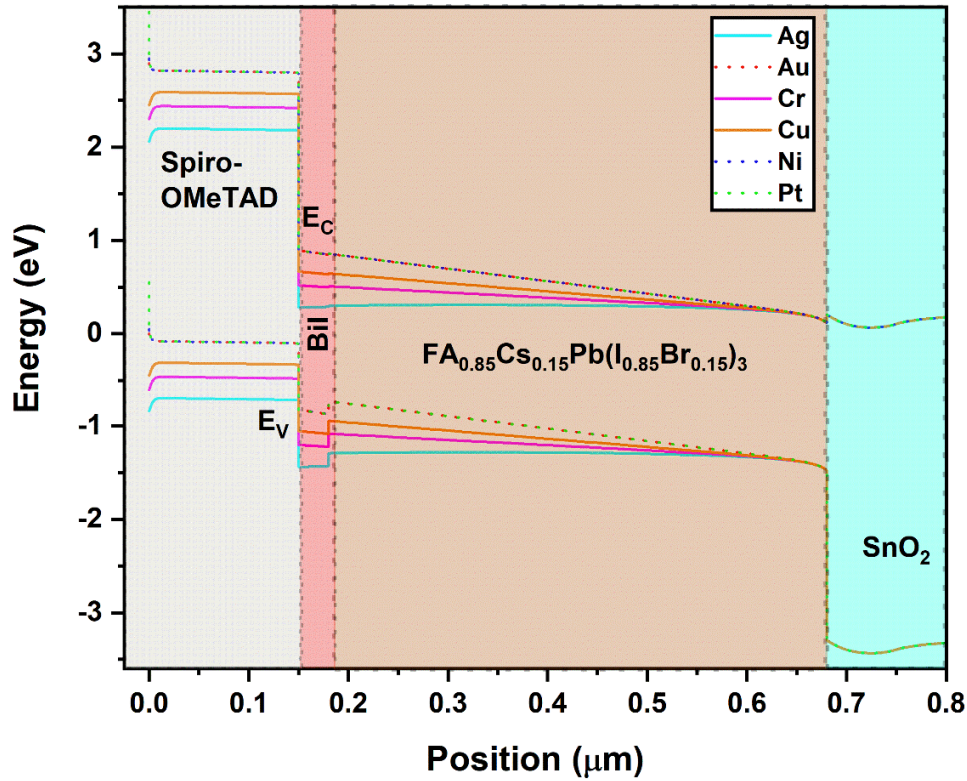
HTL		$J_{sc}$ ( $mA/cm^2$ )	FF (%)	$V_{oc}$ (V)	PCE (%)
$SrCu_2O_2$	W/o $BiI_3$	22.25	80.09	1.27	22.75
	W $BiI_3$	22.38	81.46	1.27	23.20
(Shasti and Mortezaali, 2019)		23.22	78.50	1.11	20.29
$Cu_2O$	W/o $BiI_3$	22.29	81.04	1.27	23.07
	W $BiI_3$	22.42	81.65	1.27	23.29
(Shasti and Mortezaali, 2019)		23.26	78.50	1.10	20.14
(Elseman et al., 2019)		22.53	67.36	1.13	17.23
(Lin et al., 2020)		18.41	81.73	1.06	16.03
$CuAlO_2$	W/o $BiI_3$	22.24	70.46	1.27	19.98
	W $BiI_3$	22.38	70.62	1.26	20.04
(Shasti and		23.23	78.40	1.12	20.32

<b>Mortezaali, 2019)</b>					
<b>CuSCN</b>	W/o $BiI_3$	22.25	67.05	1.27	19.05
	W $BiI_3$	22.38	66.92	1.27	19.05
<b>(Lin et al., 2020)</b>		18.28	76.34	1.08	15.18
<b>(Azri et al., 2019)</b>		21.89	83.70	1.27	23.30
<b><math>CuSbS_2</math></b>	W/o $BiI_3$	22.60	84.74	0.99	19.08
	W $BiI_3$	22.77	84.68	0.99	19.15
<b>(Devi and Mehra, 2019)</b>		31.70	81.0	0.94	24.10
<b>CuI</b>	W/o $BiI_3$	22.25	66.22	1.27	18.81
	W $BiI_3$	22.38	66.25	1.27	18.86
<b>(Lin et al., 2020)</b>		18.29	79.44	1.08	15.79
<b>(Azri et al., 2019)</b>		21.89	83.21	1.27	23.14
<b>PTAA</b>	W/o $BiI_3$	22.24	69.68	1.27	19.80
	W $BiI_3$	22.38	70.75	1.27	20.15
<b>(Coulibaly et al., 2019)</b>		41.03	74.14	0.77	23.58
<b>P3HT</b>	W/o $BiI_3$	22.26	67.10	1.20	17.99
	W $BiI_3$	22.40	66.31	1.19	17.72
<b>(Coulibaly et al., 2019)</b>		32.25	75.05	0.74	17.98
<b>(Azri et al., 2019)</b>		21.89	74.05	1.27	20.61

### 3.3 Effect of different electrodes with and without $BiI_3$ IL

In this section, the impact of various electrodes such as silver ( $Ag - 4.26 eV$ ), chromium ( $Cr - 4.5 eV$ ), copper ( $Cu - 4.65 eV$ ), gold ( $Au - 5.10 eV$ ), nickel ( $Ni - 5.15 eV$ ), platinum ( $Pt - 5.65 eV$ ) contained device performances were carefully investigated. The electrode work function (WF) values are taken from the previous reports (**Behrouznejad et al., 2016; Ming et al., 2018**). The energy band diagrams of typical and  $BiI_3$  IL included different HTL solar cells with various WF electrodes are depicted in **Fig. 7** and **Fig. S9**. The band diagram shows that the  $Au, Ni$ , and  $Pt$  lines are overlapped due to their work function value lies between 5.1 eV to 5.65 eV which is associated with their higher built in voltage behavior ( $V_{bi}$ ) (**Lin et al., 2020; Minemoto and Murata, 2014**). Other candidates such as  $Ag, Cr$ , and  $Cu$ , demonstrates the different alignment due to their lower  $V_{bi}$  and their work function values (4.2 eV to 4.65 eV) (**Lin et al., 2020**).



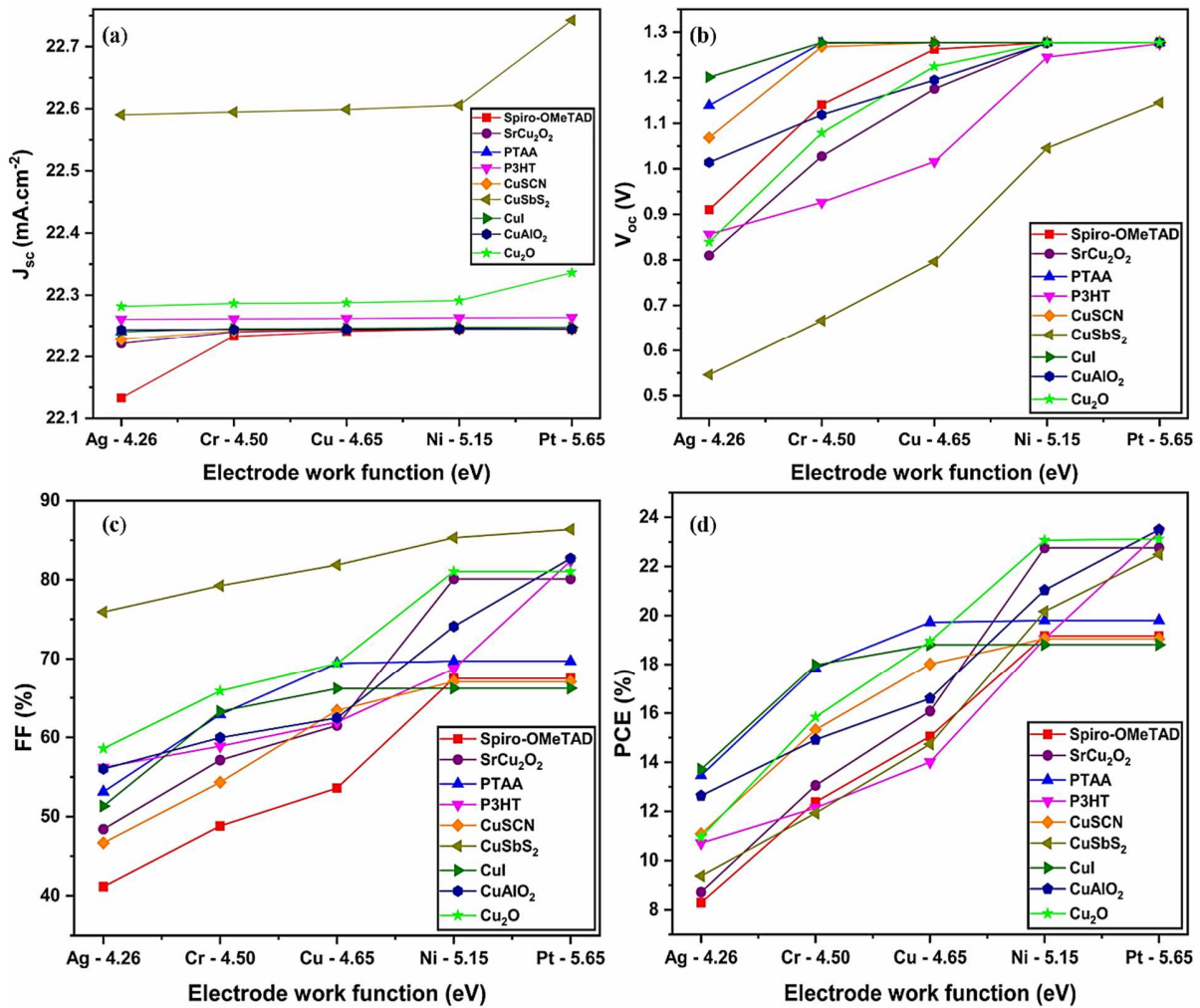


**Fig. 7** Energy band diagram for the reference device including  $BiI_3$  IL with different electrodes.

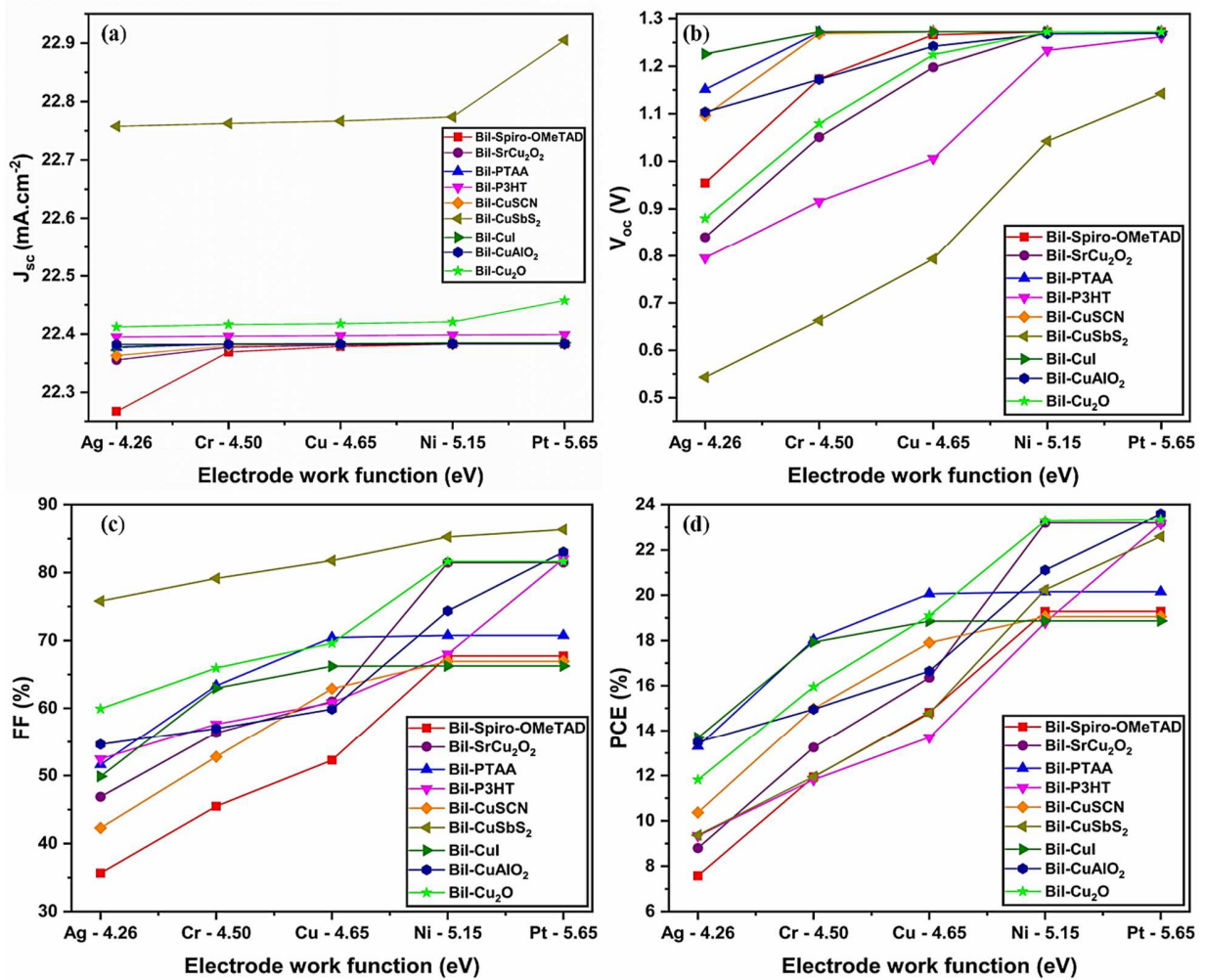
**Fig. 8** demonstrates that the obtained photovoltaic parameters ( $J_{SC}$ ,  $V_{OC}$ ,  $FF$ , and  $PCE$ ) of the typical devices with different HTL's, especially using various electrodes based on their WF values. Apart from the Perovskite/HTL interface, the energy level mismatch alignment between the HTL/electrode interface also decides the device performance (Z. Wu et al., 2019; Zheng et al., 2017). The  $CuSbS_2$ -HTL device possesses the highest  $J_{SC}$  and  $FF$  values, at the same time, it has the poor  $V_{OC}$  results than other HTL candidates. The enhanced electric field throughout the perovskite layer occurred while increasing the electrode WF value. In general, the reduced device performance corresponding to the high Schottky barriers at the HTL/electrode interface, which restricts the hole extraction (Lin et al., 2020, 2019). Most of the HTL's based devices showing similar  $PCE$  values for  $Ni$ , and  $Pt$  electrodes, except for  $CuAlO_2$ ,  $CuSbS_2$ ,  $PTAA$ , and  $P3HT$  (**Fig. 8 (d)** and Table S.II). The above mentioned HTL devices exhibits entirely enhanced photovoltaic parameters demonstrated in **Fig. 8**. It is evident that either higher or lower electrode WF directly influences the results, and also, it is related to HTL properties (i.e, hole mobility and VBO). The obtained quantum efficiency curves for typical devices are illustrated in **Fig. S10 (a)**. Simulated  $Ag$ -electrode device presenting the lower QE compared to other electrode devices, and it increases significantly based on the electrode WF values. From the QE graphs, it is easier to understand that the  $Au$ ,  $Ni$ , and  $Pt$  electrodes displaying the same behavior, the lines are overlapped which agrees the band diagram result (see **Fig. 7**). Previous reports exhibits that the oxidation from silver ( $Ag$ ) to silver iodide ( $AgI$ ) or silver bromide ( $AgBr$ ) due to the ion migration from perovskite layer through the HTL influence the device result. Also, the low



formation energy (FE) and the low diffusion barrier (DB) value (0.27 eV) of *Ag* compared to other electrodes which manipulates the device performance (Behrouznejad et al., 2016; Kato et al., 2015; Ming et al., 2018; Svanström et al., 2020). The *Cu* has much better corrosion resistance and improved DB value (0.42 eV) than *Ag*, even though it is not diffused into the perovskite layer. Therefore, *Cu* device performance was much better than *Ag* devices (Ming et al., 2018; Zhao et al., 2016). The higher DB value (0.83 eV), and the lower FE of *Cr* helps to enhance the PCE than *Ag*, meanwhile, the  $R_{Series}$  might be rising because of their high resistivity value (Ming et al., 2018). Due to the improved WF value from *Ag*, *Cr* to *Cu* electrodes (see Fig. 7), it is clear that the holes efficiently transferred from HTL to electrode, and it influence the device performances (see Fig. 8) Possibly, the higher recombination problem occurs at the perovskite/HTL or HTL/electrode interfaces in *Ag*, *Cr*, and *Cu* based devices than *Au*, *Ni*, and *Pt* devices. The higher WF electrode devices exhibit excellent photovoltaic parameters than others (see Fig. 8) (Behrouznejad et al., 2016). Therefore, it is recommended that the higher WF electrodes such as *Au*, *Ni*, and *Pt*, are more suitable candidates for mixed cation/halide perovskite solar cells.



**Fig. 8** The photovoltaic parameters ( $J_{sc}$ ,  $V_{oc}$ ,  $FF$ ,  $PCE$ ) for the reference device with different HTL's and electrodes.



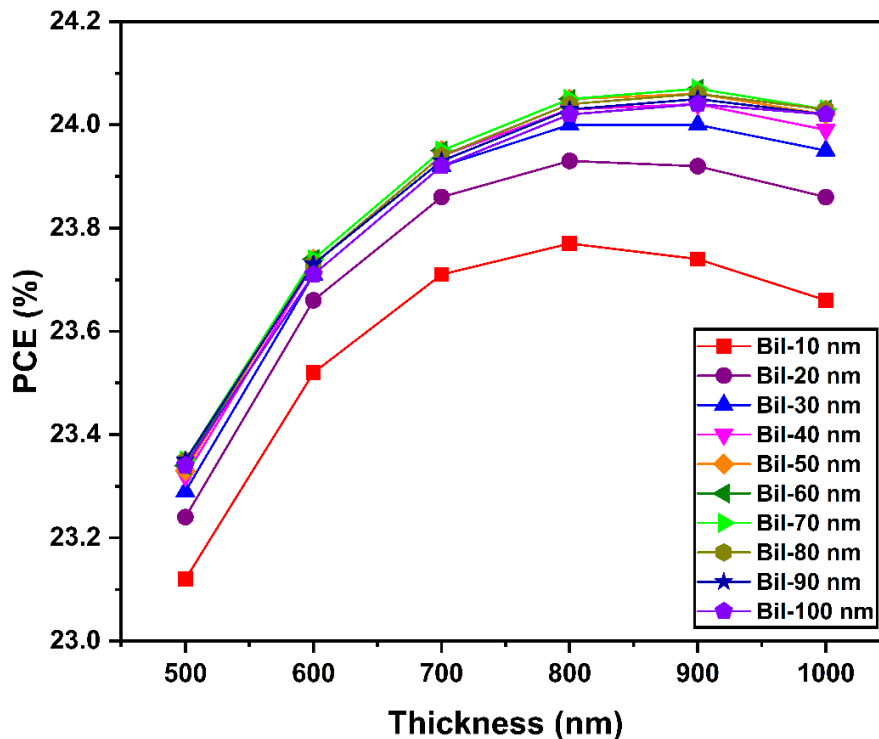
**Fig. 9** The photovoltaic parameters ( $J_{sc}$ ,  $V_{oc}$ ,  $FF$ ,  $PCE$ ) for the  $BiI_3$  IL contained devices with different HTL and electrodes.

The obtained photovoltaic parameters ( $J_{sc}$ ,  $V_{oc}$ ,  $FF$ , and  $PCE$ ) for the  $BiI_3$  IL devices with different HTL and electrodes, shown in **Fig. 9**. The obtained QE curves for the reference device, including IL, is demonstrated in **Fig. S10**. After adding the  $BiI_3$  IL, we noticed that the device performance significantly improved while using higher WF electrodes ( $Au$ ,  $Ni$ , and  $Pt$ ). Simultaneously, the *Spiro-OMeTAD*-HTL devices show a decrement in photovoltaic parameters for  $Ag$ ,  $Cr$  to  $Cu$  electrodes due to their lower WF value as well as it possibly creates the high Schottky barrier at the HTL/electrode interface which reduce the hole transport from HTL to an electrode. (Table S.II) The case of *P3HT*-HTL devices (for all anodes) demonstrates a slightly lowered performance than the normal one. Most HTL candidates prove that a thin  $BiI_3$  IL efficiently enhances the device results with all the electrodes mentioned above. The device

performance is strongly affected due to the formation of the resistance layer at the HTL/electrode interface, so it is crucial to use the appropriate valence and conduction band energy levels having HTL (Behrouznejad et al., 2016). Moreover, the reported experimental result shows that placing a thin *Cr* interfacial layer inbetween HTL and *Au* interface significantly reduced the electrode diffusion at the higher temperature, which is beneficial for higher performance devices (Domanski et al., 2016). Therefore, adding an IL between the interfaces could be an excellent strategy to improve solar cell performance. In the following section, we will demonstrate the influence of *BiI<sub>3</sub>* IL thickness over the absorber layer thickness.

### 3.4 Effect of *BiI<sub>3</sub>* IL and absorber layer thickness

Generally, the properties of the perovskite absorber layer determine the device's performance. Mainly, the layer thickness is one of the leading parameter to enhance photon absorption, significantly improving the device results. After careful investigation, the *Cu<sub>2</sub>O*-HTL was selected to perform the mentioned thickness analysis. Therefore, in this section, the influence of the perovskite layer and the *BiI<sub>3</sub>* IL thickness over the *Cu<sub>2</sub>O*-HTL based device performance was carefully investigated. The absorber layer thickness differs from 500 to 1000 nm, as well as the interfacial layer thickness varied from 10 to 100 nm, and the other listed parameters in Table I remain the same.

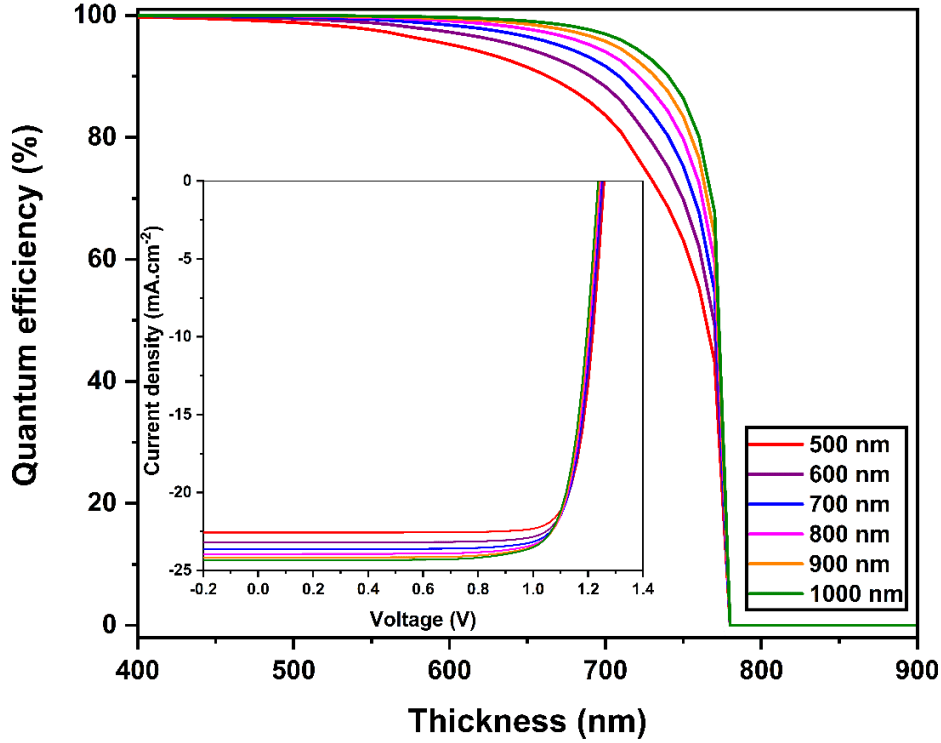


**Fig. 10** The variation of PCE (a) with different perovskite layer (500 nm to 1000 nm) and *BiI<sub>3</sub>* IL (10 nm to 100 nm) thickness.

**Fig. 10** demonstrates that the device PCE considerably enhances (23 % to 24 %) while increasing the absorber layer thickness from 500 to 1000 nm. It continuously improves up to 900 nm, then

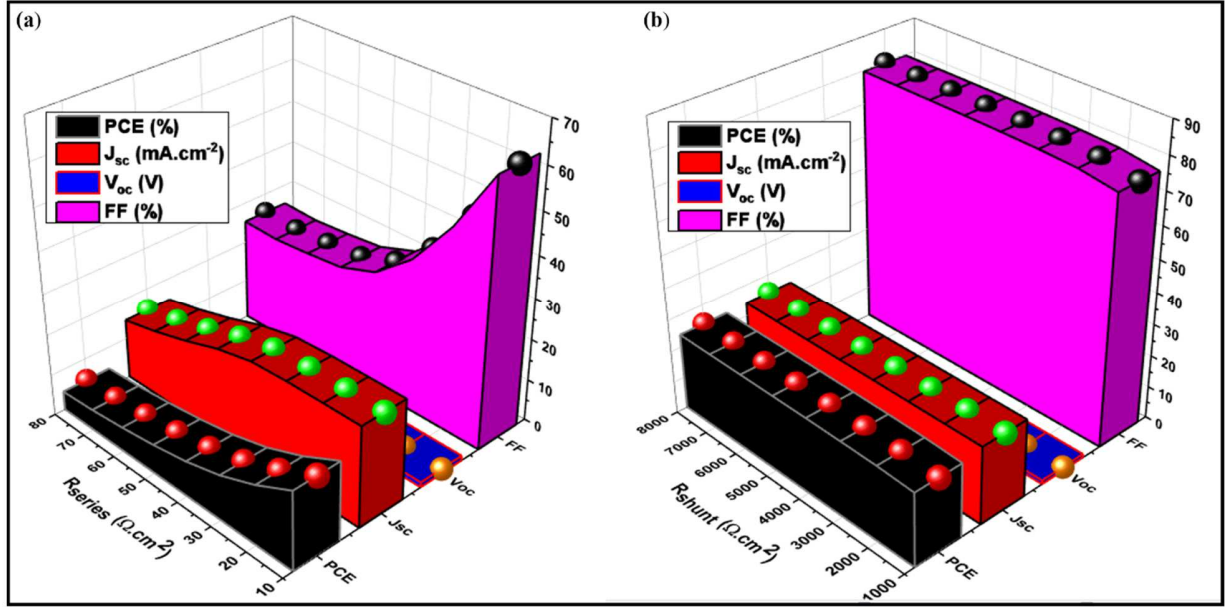
slightly decreases. Several experimental reports reveal that a thick perovskite layer will absorb more photons, especially in the red to the near-infrared region (Chen et al., 2017; Yuan et al., 2016), which is appropriate for solar devices. For example, Jiehuan Chen et al. studied the solar cell performance based on the different perovskite layer thickness, especially beyond 700 nm (i.e., 700, 850, 1000, 1150, 1400, and 1650 nm) using a simple hot casting method (Chen et al., 2018). The 850 nm thicker absorber perovskite device provides the highest PCE (19.54%) among all other absorbers. Noticeably, the device maintains the stable PCE value of 19% from 700 to 1150 nm thickness range, and the device holds 80% of its initial PCE after 30 days. Therefore, using an optimum thicker perovskite absorber and the IL is beneficial for device improvement (Sanehira et al., 2016; S. Wu et al., 2019). Shaohang Wu et al. studied the inverted perovskite solar cell PV performance and stability while inserting a bismuth (Bi) interlayer (thickness - 5, 10, 20, 40, and 80 nm) (S. Wu et al., 2019). The results show that the thicker Bi film offers a uniform and good coverage film than a thin layer. Due to increased contact resistance, a thicker Bi layered (i.e., 40 nm) solar cell reduces the device efficiency; most importantly, thicker IL lessens the degradation induced from ambient moisture leading to the stability enhancement (6000 h).

In that perspective, we simultaneously investigated the  $BiI_3$  IL thickness with different perovskite layer thickness and makes it evident that the increment substantially improves the PCE (see **Fig. 10**). So, the optimum  $BiI_3$  IL thickness was found at 70 nm, which corresponds to the superior PCE (i.e.,  $PCE = 24.07\%$ ,  $J_{sc} = 24.18 \text{ mA/cm}^{-2}$ ,  $FF = 80.25\%$ , and  $V_{oc} = 1.24 \text{ V}$ ). Moreover, the thick perovskite (more than 900 nm) layered device possibly suffered from higher recombination issues at the interface due to the absorber thickness that is larger than the diffusion length of charge carriers, which led to poor device performance. **Fig. 11** exhibits that the obtained QE graph's tendency is consistent with the PCE results (see **Fig. 10**). The current density – voltage characteristics (inset) also proves the previously mentioned behavior.



**Fig. 11** QE and J-V characteristics (inset) of a different perovskite absorber layer thickness (500 nm to 1000 nm), including 70 nm of  $BiI_3$  IL.

The simulation findings suggest that the 900 nm perovskite layer with 70 nm  $BiI_3$  IL is appropriate for an excellent FAPI based perovskite device. However, the parasitic resistances ( $R_{series}$  and  $R_{shunt}$ ) are largely influence the experimental device performance. For example, a higher efficiency device must be having a low  $R_{series}$  and a high  $R_{shunt}$ . Our previous work clearly demonstrates that the effect of these resistances over three different perovskite absorber layer based devices (Karthick et al., 2020). Therefore, here we investigated the impact of these resistances over the  $Cu_2O$ -HTL device to understand the device behavior. We varied the  $R_{series}$  and  $R_{shunt}$  values between 10 – 80  $\Omega.cm^2$  and 1000 – 5000  $\Omega.cm^2$ , respectively, and the corresponding evolution of the photovoltaic parameters are demonstrated in Fig. 12 (a) and (b). Our obtained results show that an increase in  $R_{series}$  from 10 to 80  $\Omega.cm^2$  strongly decreases the FF from 63% to 25%. Moreover, for a  $R_{series}$  over 40  $\Omega.cm^2$ ,  $J_{sc}$  is rapidly reduced from 24 to 14  $mA/cm^2$  and the overall PCE is rapidly decreased from 24% to only 4.5% (see Fig.12 (a)). Usually,  $R_{series}$  does not influence the  $V_{oc}$ , and the obtained  $V_{oc}$  behavior is shown in Fig.12 (a). In comparison, increasing  $R_{shunt}$  enhances the FF and PCE values (see Fig.12 (b)), at the same time, it's slightly affects the  $V_{oc}$  and  $J_{sc}$  values. So, the high  $R_{shunt}$  is a indication of less defect and interface states presence in the device. While varying the  $R_{series}$  and  $R_{shunt}$  values into the optimized device considerable changes occur in the photovoltaic parameters. Therefore, it is clear that these parasitic resistances efficiently affect device performance.



**Fig. 12** Effect of  $R_{series}$  (a) and  $R_{shunt}$  (b) in the performance of simulated  $Cu_2O$ -HTL device including  $BiI_3$  IL.

Finally, we also employed the optimized thicknesses (absorber and IL layer) in  $SrCu_2O_2$ -HTL solar cell and it shows a PCE of 23.91 % with a short current density ( $J_{sc}$ ) value of  $24.17 \text{ mA/cm}^{-2}$ , a fill factor (FF) of 79.73 %, and an open-circuit voltage ( $V_{oc}$ ) of 1.24 V.

In the end, this paper does not address any recombination issues (such as first, second, and third-order), which generally occurs in the real perovskite device. From the experimental point of view, the following points are crucial for boosting the actual perovskite device performance, such as (1) enhancing the absorber layer phase and its morphology, (2) improving the charge transport behavior, (3) passivating the surface nor interface defects, (4) altering the energy level alignment between the ETL/perovskite and perovskite/HTL interfaces, (5) improving the stability, (6) reducing the hysteresis. In general, to remove the traps at the interfaces is beneficial for the charge transfer. In that case, an interlayer's addition to the device helps to minimize the traps and suppress the ion immigration. Therefore, interface modifications and the chosen interfacial material directly influences the device stability, performance, and hysteresis behavior. Adding polymers, small molecules, 2D materials, or self-assembled monolayers (SAM) in between the perovskite/HTL interface efficiently may improve the charge transfer from the absorber to HTL and electrode (Shao and Loi, 2020). As a result, based on our report, a further insightful experimental study is necessary to understand the mechanism behind the  $BiI_3$  IL over the device performance.

#### 4. Conclusion

In this work, for the first time  $FA_{0.85}Cs_{0.15}Pb(I_{0.85}Br_{0.15})_3$  perovskite absorber layer based n-i-p planar device with and without  $BiI_3$  interfacial layer tested with different HTL's and various electrodes, which was carefully analyzed and presented using SCAPS. The simulated results exhibit that the  $Cu$  based HTL's, especially  $Cu_2O$  &  $SrCu_2O_2$  shows excellent photovoltaic



performance than other HTL competitors. The optimized 900 nm perovskite layer with a thin (70 nm) interfacial  $BiI_3$  layer has a more substantial influence on the solar cell results, which significantly boosted the photovoltaic parameters by reducing the perovskite/HTL interface defects. Therefore, both  $Cu_2O$  and  $SrCu_2O_2$  HTL's are considered as suitable HTL's (with and without  $BiI_3$  IL) for  $FA_{0.85}Cs_{0.15}Pb(I_{0.85}Br_{0.15})_3$  based planar device. Also, the impact of different electrodes was examined using various HTL candidates, and our results emphasize that the higher work function electrodes such as  $Au$ ,  $Ni$ , and  $Pt$  are more considerable than  $Ag$ ,  $Cr$  to  $Cu$ . From the result, the simulated device displaying an efficiency of around 24 % with a suitable  $Cu_2O$ -HTL. Here, the absence of parasitic electrical losses (without considering the optical reflectance of each layer at the interface nor surface) could be the prominent reason for the obtained higher efficiency than the reported experimental reports. The optimized absorber, IL thickness, and parasitic resistances are also crucial for efficient devices. These simulation findings are useful for further understanding of the  $BiI_3$  IL in between perovskite/HTL interface and enhancing the mixed cation/halide perovskite solar cells.

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